1.3.95

Dear Abner,

I was very sorry we were not able to continue our discussions, after the unfortunate mishap with my back. It is slowly improving, but I am still not able to give my lectures. However Pagonis and Dickson have nobly stood in for me, and their lectures seem proving quite a bit more popular with the students than my own!

It was very nice to see you again, and I hope you enjoyed your visit as much as we enjoyed having you.

I read through your 'theorem' that Jeremy passed on to me. I agree with everything you say. My condition $* \quad \forall P_2 \exists P_1 \text{ s.t. } \langle P_1 P_2 \rangle_{\Psi} = \langle P_1 \rangle_{\Psi}$ is true iff Ψ is entangled, and the baby Reeh-Schlieder theorem follows iff Ψ is entangled. So $*$ is necessary and sufficient for $P.\Psi$

the Baby R-S theorem.

This is what I claimed in 'More Ado--' - sufficiency on p-12, and necessity as Theorem 4 on p. 14.

The reason I proved Theorem 4 in such an apparently roundabout way was the following.

I could have argued simply that if * were false and hence $\mathcal{H}$ not entangled, then the baby R-S theorem would be false. So, contrapositively Baby R-S $\Rightarrow$ *.

But I wanted to give a 'direct' proof of *, $\underline{i.e.}$ using only Baby R-S, which could then be lifted straight back to the field theory case, which was the one I was really interested in. This is, heuristically, how I hit on the idea that the vacuum correlations in the field theory case are maximal in the sense of * (modulo suitable epsilonics!)

My discussion of Baby R-S in 'More Ado--' was really meant to explain to the reader how I came to think that the main result of the paper ('Theorem 4' on p. 18) was a plausible thing to try and prove!

with very best wishes
Michael